Waves in unbounded photonic crystals and transmission properties at interfaces

Ben Schweizer

technische universität
dortmund

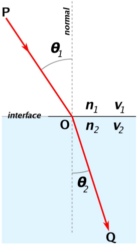
August 8, 2017

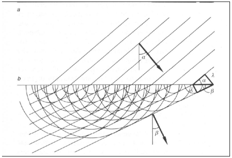
Geometric optics vs. Wave equation

Fermat's principle of **the fastest path:**

Light finds the fastest way to reach the destination!

$$\frac{\sin\Theta_1}{\sin\Theta_2} = \frac{v_1}{v_2} = \frac{n_2}{n_1}$$

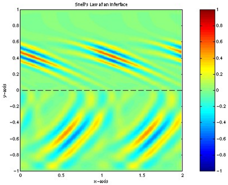




Huygens' principle of **superpositions**

Wave equation

$$\partial_t^2 u = \Delta u$$



Numerical solution

Maxwell's equations and negatve index

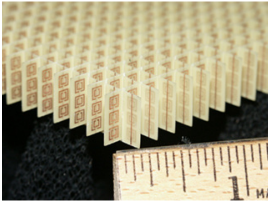
Maxwell's Equations (1865)

$$\begin{aligned} \text{curl } E &= i\omega\mu H \\ \text{curl } H &= -i\omega\varepsilon E \end{aligned}$$

E: electric field, H: magnetic field

$$H, E \sim e^{-i\omega t}$$

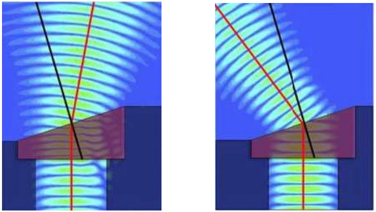
- $\mathrm{Re}\,\varepsilon < 0$ possible
- μ is always 1
- $\mathrm{Re}\,\mu\varepsilon < 0$: medium is opaque



Veselago (1968)

Materials with negative index

$\varepsilon < 0$ and $\mu < 0 \Rightarrow$ negative index!

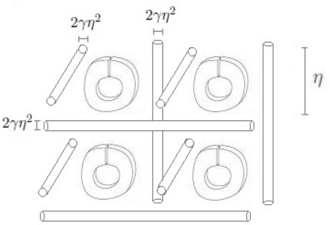


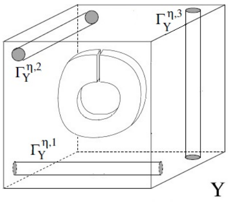
Solutions for positive and negative index

Pendry et al. ($\sim$ 2000)

Design of a negative index meta-material
Use split rings and wires

Microscopic geometry with wires





(H^η, E^η) solves Maxwell, $(H^\eta, E^\eta) \to (\hat{H}, \hat{E})$ *"geometrically"*

Effective Maxwell system (A.Lamacz & B.S., SIAM J.Math.Anal. 2017)

$$\operatorname{curl} \hat{E} = i\omega\mu_{\text{eff}}\, \hat{H}$$
$$\operatorname{curl} \hat{H} = -i\omega\varepsilon_{\text{eff}}\, \hat{E}$$

with negative (for appropriate geometry and $\mathrm{Re}(\varepsilon_w) < 0$) coefficients

$$\mu_{\text{eff}} = \mu_{\text{eff,R}} \quad \text{and} \quad \varepsilon_{\text{eff}} = \varepsilon_{\text{eff,R}} + \pi\gamma^2\, \varepsilon_W .$$

Wave transmission into photonic crystals

Our motivation:

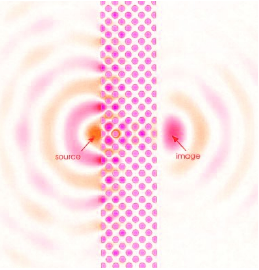
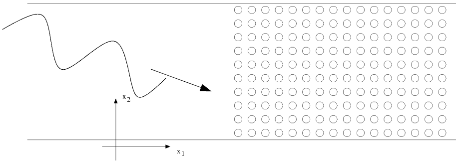


Image taken from

C. Luo, S. G. Johnson, J. D. Joannopoulos, and J. B. Pendry. **All-angle negative refraction without negative effective index.** *Phys. Rev. B, 65:201104, May 2002*

Is this negative refraction at a photonic crystal?



Geometry of the transmission problem.

We study the waves that are generated in the photonic crystal.

Helmholtz equation:

$$-\nabla \cdot (a \nabla u) = \omega^2 u$$

Bloch expansion: Write an arbitrary function f in a smart way!

1.) $f : \mathbb{R}^n \to \mathbb{R}$ is written with a Fourier transform:

$$f(x) = \int_{\mathbb{R}^n} \hat{f}(\xi) e^{2\pi i \xi \cdot x} \, d\xi$$

2.) ξ is written as $\xi = k + j$ with $k \in \mathbb{Z}^n$ and $j \in [0,1)^n =: Z$

$$f(x) = \int_Z \underbrace{\sum_k \hat{f}(k+j) e^{2\pi i k \cdot x}}_{=:F} e^{2\pi i j \cdot x} \, dj$$

3.) Periodic $F = F(x; j)$ is expanded in periodic eigenfunctions $\Psi_{j,m}(x)$:

$$F(x; j) = \sum_{m \in \mathbb{N}} \alpha_{j,m} \Psi_{j,m}(x)$$

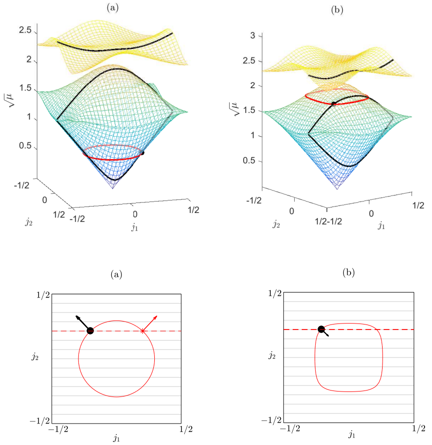
$U_{j,m}(x) := \Psi_{j,m}(x) e^{2\pi i j \cdot x}$ solves

$$-\nabla \cdot (a(x) \nabla U_{j,m}(x)) = \mu_{j,m} U_{j,m}(x)$$

Result: The operator $L = -\nabla \cdot (a(.)\nabla)$ acts as a multiplier:

$$f(x) = \int_Z \sum_{m \in \mathbb{N}} \alpha_{j,m} U_{j,m}(x) \, dj, \qquad Lf = \int_Z \sum_{m \in \mathbb{N}} \alpha_{j,m} \mu_{j,m} U_{j,m}(x) \, dj$$

Energy landscape in the periodic medium



The three surfaces correspond to $m = 0, 1, 2$.
The vertical axis shows $\sqrt{\mu_{j,m}}$, where $\mu_{j,m}$ is the m-th eigenvalue for the wave vector $j = (j_1, j_2)$.

The arrows show gradients of the energy landscape

Energy transport by Bloch waves

For $\lambda = (j, m)$, which Bloch wave U_λ is "right going"?

Recall for Maxwell: The Poynting vector $P := E \times H$ measures the energy flux

Poynting number

The Poynting number P_λ describes the right-going energy:

$$P_\lambda := \operatorname{Im} \fint_{Y_\varepsilon} \bar{U}_\lambda(x)\, e_1 \cdot [a(x) \nabla U_\lambda(x)]\; dx$$

Index sets: Left-going waves and "vertical waves"

$$I_{<0} := \{\lambda \in I \mid P_\lambda < 0\} \qquad \text{and} \qquad I_{=0} := \{\lambda \in I \mid P_\lambda = 0\}$$

Projection: Onto left-going waves

$$\Pi_{<0} u(x) := \sum_{\lambda \in I_{<0}} \alpha_\lambda U_\lambda(x)$$

Radiation for homogeneous media: Sommerfeld, 1912

Homogeneous problem $-\Delta u = \omega^2 u$ in $\mathbb{R}^n$

Fundamental solutions

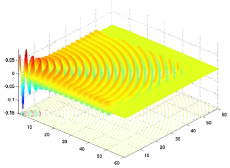
Two fundamental Helmholtz solutions for $x \in \mathbb{R}^3$:

$$u_+(x) := \frac{1}{|x|} e^{i\omega|x|} \quad \text{and} \quad u_-(x) := \frac{1}{|x|} e^{-i\omega|x|}$$

Time-dependence $e^{-i\omega t}$ implies: u_+ is an outgoing wave, u_- an incoming wave.

Sommerfeld condition

$$|x|^{(n-1)/2} (\partial_{|x|} u - i\omega u)(x) \to 0 \quad \text{as } |x| \to \infty$$

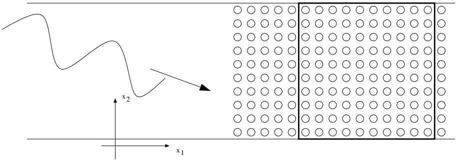


Both elementary solutions decay for $|x| \to \infty$. It is not reasonable to demand only a decay property

- u_+ satisfies the Sommerfeld condition
- u_- does not

Justification (Sommerfeld): Radiation condition implies uniqueness

Expansion of solutions



We consider u only on the marked square

After a shift:
$u \in L^2((0, R\varepsilon) \times (0, R\varepsilon))$

Wave-vector: $j \in Z := [0,1)^2$. Eigenvalue number: $m \in \mathbb{N}_0$
Multiindex: $\lambda = (j, m) \in I_K$. Basis: $U_\lambda^+(x) := \Psi_\lambda^+(x)\, e^{2\pi i \theta(\lambda)\cdot x/\varepsilon}$

$$u(x) = \sum_{\lambda \in I_K} \alpha_\lambda^+ U_\lambda^+(x)$$

Expansion of an arbitrary function u in Bloch waves

For "outgoing solutions" we demand (on the right):

u consists only of right-going Bloch modes

S. Fliss and P. Joly. Solutions of the time-harmonic wave equation in periodic waveguides: asymptotic behaviour and radiation condition. Arch. Ration. Mech. Anal., 219, 2016

Our wish-list

Transmission problem

a constant on the left, periodic on the right

Helmholtz equation: $-\nabla \cdot (a \nabla u) = \omega^2 u$, periodic in vertical direction

Outgoing wave conditions, on the right:

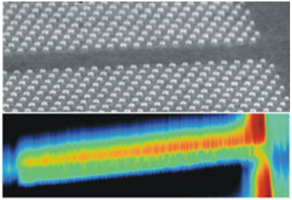
$$\fint_{RY_\varepsilon} \left| \Pi^+_{<0}(u^+_R) \right|^2 \to 0 \quad \text{as} \quad R \to \infty$$

Wishful thinking: For *every* frequency $\omega > 0$

- There exists a solution to the problem
- The solution to the problem is unique

Uniqueness cannot be expected

There are surface-waves $\longrightarrow$ no uniqueness!



S. Bozhevolnyi/Aalborg Univ.

Bloch measures

G. Allaire and C. Conca. Bloch wave homogenization and spectral asymptotic analysis. J. Math. Pures Appl. 1998

Let $u_R \in L^2(W_R;\mathbb{C})$ be a sequence

$$u_R(x) = \sum_{\lambda\in I_R} \alpha_\lambda^\pm U_\lambda^\pm(x)$$

Discrete Bloch-measure for fixed $l \in \mathbb{N}_0$:

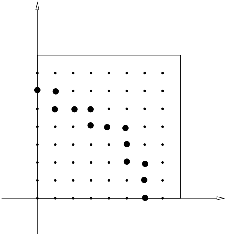
$$\nu_{l,R}^\pm := \sum_{\lambda=(j,l)\in I_R} |\alpha_\lambda^\pm|^2 \;\; \delta_j$$

where δ_j denotes the Dirac measure in $j \in Z$.
If, as $R \to \infty$,

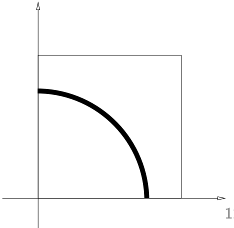
$$\nu_{l,R}^\pm \to \nu_{l,\infty}^\pm$$

in the sense of measures, then

$\nu_{l,\infty}^\pm \in \mathcal{M}(Z)$ is a Bloch measure generated by u



The Brillouin zone $Z = [0,1)^2$.
A periodic u is expanded with discrete values of $j \in Z$.



Uniqueness result

Frequency assumption with Bloch-eigenvalues $\mu_m^{\pm}(j)$:

$$\omega^2 < \inf_{j \in Z,\, m \geq 1} \mu_m^{+}(j)$$

Theorem (A.Lamacz & B.S., Uniqueness)

Let u and $\tilde{u}$ be two solutions of the transmission problem. Then the difference $v := u - \tilde{u}$ generates a Bloch measure that has support only on vertical waves.

Interpretation: Waves can be

- localized at the interface

  or

- travelling vertically in the photonic crystal

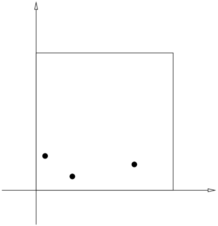
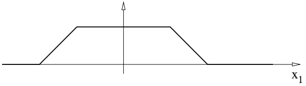


Figure: The indices $j \in Z$ corresponding to "vertical waves"

Uniqueness follows from energy conservation

Let v solve the Helmholtz equation with coefficients $a = a^\varepsilon$. Use

$$\vartheta(x) := \begin{cases} 1 & \text{if } |x_1| \le \varepsilon R \\ 2 - \frac{|x_1|}{\varepsilon R} & \text{if } \varepsilon R < |x_1| < 2\varepsilon R \\ 0 & \text{if } |x_1| \ge 2\varepsilon R \end{cases}$$



and the test-function $\vartheta(x)\,\overline{v}(x)$ to obtain

$$\int_{\mathbb{R}} \int_0^h \left\{ a^\varepsilon\, \vartheta\, |\nabla v|^2 + a^\varepsilon\, \partial_{x_1}\vartheta\ \overline{v}\ \partial_{x_1} v \right\} = \omega^2 \int_{\mathbb{R}} \int_0^h \vartheta\, |v|^2$$

Poynting vector bilinear form $b_R^\pm : L^2(W_R;\mathbb{C}) \times H^1(W_R;\mathbb{C}) \to \mathbb{C}$:

$$b_R^+(u,v) := \int_{W_R} \bar{u}(x)\, e_1 \cdot [a^\varepsilon(x) \nabla v(x)]\, dx$$

Take the imaginary parts and obtain the energy conservation

$$\operatorname{Im} b_R^- \left(v_R^-, v_R^-\right) = \operatorname{Im} b_R^+ \left(v_R^+, v_R^+\right)$$

Result: If both terms have opposite sign, they must vanish!

Show $\nu^{\pm}_{l,\infty} = 0$ for $l \geq 1$

Let $\delta > 0$ be a number with $\delta \leq |\omega^2 - \mu_{(j,m)}|^2$ for all j and $m \geq 1$. Then, formally,

$$\begin{aligned}
\delta \int_{W_R} \left|\Pi^{\mathrm{ev},+}_{m\geq 1}(u^+_R)\right|^2 &= \delta \sum_{\substack{\lambda=(j,m)\in I_R \\ m\geq 1}} \left|\left\langle u^+_R , U_\lambda\right\rangle_R\right|^2 \\
&\leq \sum_{\substack{\lambda=(j,m)\in I_R \\ m\geq 1}} \left|(\omega^2 - \mu_\lambda)\langle u^+_R , U_\lambda\rangle_R\right|^2 \\
&\leq \sum_{\lambda\in I_R} \left|\left\langle \omega^2 u^+_R , U_\lambda\right\rangle_R - \left\langle \mu_\lambda u^+_R , U_\lambda\right\rangle_R\right|^2 \\
&= \sum_{\lambda\in I_R} \left|\left\langle \mathcal{L}_0(u^+_R) , U_\lambda\right\rangle_R - \left\langle \mu_\lambda u^+_R , U_\lambda\right\rangle_R\right|^2 = 0
\end{aligned}$$

The calculation can be made precise with cut-off functions on large squares. **Result for Bloch measure:** $\nu^{\pm}_{l,\infty} = 0$ for $l \geq 1$

A similar calculation yields: $\mathrm{supp}(\nu^{\pm}_{0,\infty}) \subset \left\{ j \in Z \,|\, \mu^{\pm}_0(j) = \omega^2 \right\}$

Transmission conditions

Assume again: Frequency below second band
The vertical wave number is conserved:

Theorem (Transmission conditions)

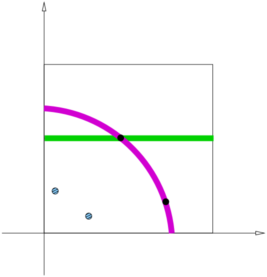
Let u be a solution of the transmission problem.
Let $\nu^{\pm}_{l,\infty}$ be a Bloch measure to u.

Then: $\nu^{\pm}_{l,\infty} = 0$ *for* $l \geq 1$,

$$\mathrm{supp}(\nu^{\pm}_{0,\infty}) \ \subset \ \{j \in Z \,|\, \mu^{\pm}_0(j) = \omega^2\}$$

and

$$\mathrm{supp}(\nu^{\pm}_{0,\infty}) \ \subset \ \{j \in Z \,|\, j_2 = k_2\} \ \cup \ J^{\pm}_{=0,0}$$



Waves must have:

- the correct energy **and**
- the correct k_2 (or be vertical)

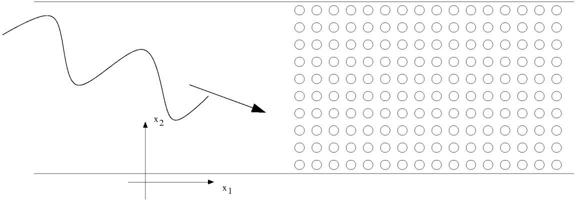
The theorem follows from uniqueness: Compare u with its projection to the vertical wave number k_2

A numerical scheme

Based on the radiation condition $\longrightarrow$ numerical scheme

T. Dohnal and B. Schweizer: A Bloch wave numerical scheme for scattering problems in periodic wave-guides (submitted)

$$-\nabla \cdot (a\nabla u) = \omega^2(1 + \mathrm{i}\delta)u + f$$



Concept:

- At the far left/right:
  Solution is a linear combination of outgoing Bloch waves
- Standard finite elements in the core domain

Design of the scheme, negative refraction

In the radiation boxes $W^{\pm}_{R,L}$ use $X^{\pm} := span\{U^{\pm}_{\lambda} \,|\, \lambda \in I^{\pm}\}$,
The index sets $I^{\pm}$ satisfy $\lambda \in I^{\pm} \Rightarrow \pm P^{\pm}_{\lambda} > 0$.

Function space:

$$V := \left\{ u \in H^1(\Omega_{R+L}) \,\middle|\, u \text{ vertically periodic, } \{u\}^{+}_{R,L} \in X^{+}, \{u\}^{-}_{R,L} \in X^{-} \right\}$$

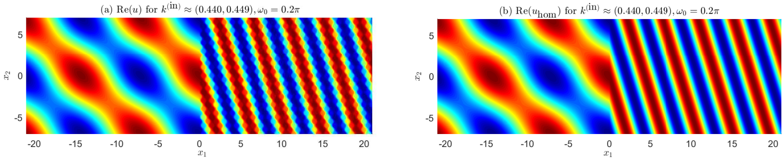
Bilinear form (with cut-off function ϑ as above):

$$\beta(u,v) := \int_{\Omega_{R+L}} a\nabla\bar{u}\cdot\nabla v\,\vartheta - \int_{\Omega_{R+L}} (1 - \mathrm{i}\delta\mathbf{1}_{\Omega_R})\,\omega^2\bar{u}\,v\,\vartheta - \frac{1}{\varepsilon L}\int_{W^{+}_{R,L}} a\nabla\bar{u}\cdot e_1\,v + \frac{1}{\varepsilon L}\int_{W^{-}_{R,L}} a\nabla\bar{u}\cdot e_1\,v = \int_{\Omega_R} \bar{f}\,v$$

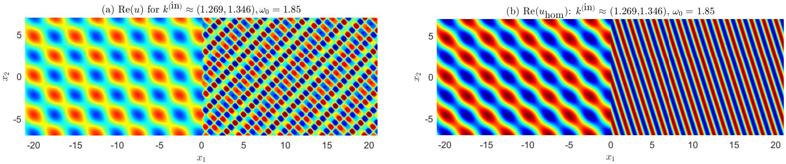
Coercivity of β follows from $\nabla\vartheta = \mp\frac{1}{L}e_1$ and $P^{\pm}_{\lambda} > 0$.

Numerical results: Comparison with homogenization

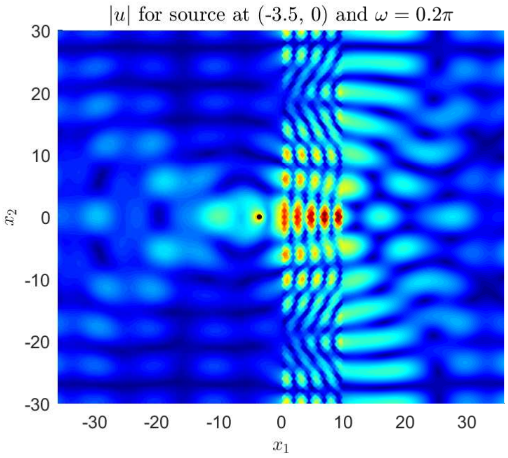
Transmission into periodic medium I: Large wave-length



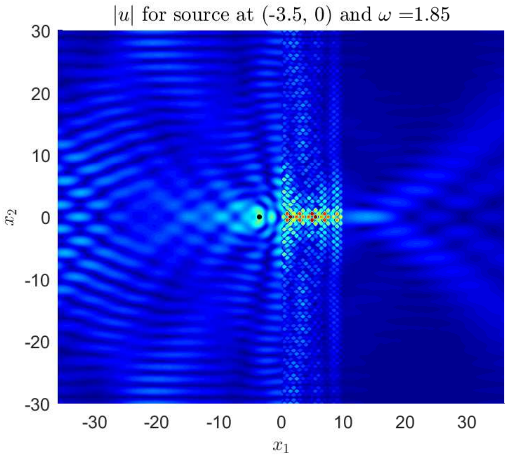
Transmission II: Wave-length comparable to structure



A finite crystal with positive refraction property



A finite crystal with negative refraction property



Thank you!