

**Energy method approach to existence  
results for the Helmholtz equation in  
periodic wave-guides**

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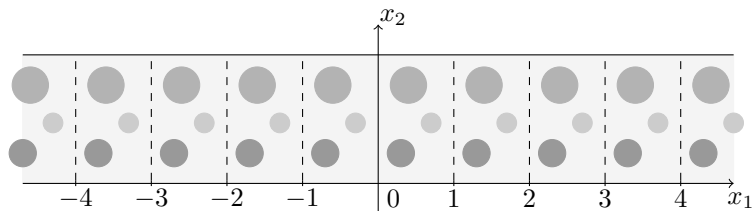
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# Helmholtz equation in wave guide

Wave-guide geometry:  $\Omega := \mathbb{R} \times S$  with  $S \subset \mathbb{R}^{d-1}$  bounded Lipschitz

Coefficient:  $a : \Omega \rightarrow \mathbb{R}$  strictly positive, 1-periodic in  $x_1$

Neumann condition on  $\partial\Omega$



Given data:  $f \in L^2(\Omega)$  with compact support, frequency  $\omega \in \mathbb{R}$

Helmholtz equation

$$-\nabla \cdot (a \nabla u) = \omega^2 u + f \quad \text{in } \Omega \quad (\text{H})$$

**Main result: existence of solutions**

Important is the method: Only energy methods!

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# Main result on periodic wave-guides

**Flux conservation:** For solutions  $\varphi$  of (H) and  $-\infty < \rho < r < \infty$  holds

$$\operatorname{Im} \int_{\{\rho\} \times S} a \nabla \varphi \cdot e_1 \bar{\varphi} = \operatorname{Im} \int_{\{r\} \times S} a \nabla \varphi \cdot e_1 \bar{\varphi}$$

**Solution concept:**  $u \in H_{\text{loc}}^1(\Omega)$  with

- (i)  $u$  solves (H) in  $\Omega$  in the weak sense
- (ii)  $\sup_{r \in \mathbb{Z}} \|u\|_{L^2((r, r+1) \times S)} < \infty$
- (iii) the radiation condition is satisfied

**A loose formulation of (iii) on the right:** There exist finitely many quasiperiodic homogeneous solutions  $\varphi_j$  of (H) with positive energy-flux such that  $u - \sum_j \alpha_j \varphi_j \rightarrow 0$  as  $x_1 \rightarrow \infty$ .

**Theorem (Existence and uniqueness result for periodic media)**

*Let data  $\Omega$ ,  $f$ ,  $\omega$ , and  $a$  be fixed,  $a(x + e_1) = a(x) \forall x \in \Omega$ . Let  $\omega$  be a non-singular frequency (see below). Then there exists one and only one solution  $u$  to the radiation problem (i)–(iii).*

# Non-singular frequency

We use  $Au := -\nabla \cdot (a\nabla u)$ , cylinders  $W_r := (r, r+1) \times S$ , the norm  $\|u\|_{sL} := \sup_{r \in \mathbb{Z}} \|u\|_{L^2(W_r)}$ , and the space

$$X := \{u|_{W_0} \mid u \in H_{\text{loc}}^1(\Omega), \|u\|_{sL} < \infty, Au = \omega^2 u \text{ in } \Omega\}$$

## Definition (Non-singular frequency)

$\omega > 0$  is called a non-singular frequency (for coefficient  $a$ ) if:

- (a) Finite dimension: The space  $X$  has a basis  $(\varphi_j)_{1 \leq j \leq M}$  with quasimoments  $\xi_j \in [0, 2\pi)$  such that each  $\varphi_j$  possesses a  $\xi_j$ -quasiperiodic extension satisfying  $A\varphi_j = \omega^2 \varphi_j$  in  $\Omega$ .
- (b) Non-vanishing flux: For every quasiperiodic function  $u \in H_{\text{loc}}^1(\Omega)$  with  $Au = \omega^2 u$ , the restriction  $\varphi = u|_{W_0} \in X$  has the property

$$\text{Im} \int_{W_0} a \nabla \varphi \cdot e_1 \bar{\varphi} \neq 0.$$

Exclusion of singular frequencies appears in all “blue” references

# Right-going waves, projections, radiation condition

Construct basis  $\phi_j^\pm$ , flux of  $\phi_j^+/\phi_j^-$  is positive/negative

$u \in X$  can be written as  $u = \sum_{j=1}^N \alpha_j \phi_j^+ + \sum_{j=1}^N \beta_j \phi_j^-$

With, e.g.,  $\Pi_{X,+} : u \mapsto \sum_{j=1}^N \alpha_j \phi_j^+$  and orthogonal projection:

$$\Pi_\pm : L^2(W_0) \rightarrow X_\pm \subset L^2(W_0)$$

$\Pi_+/\Pi_-$  projects onto right/left going base waves

**Definition (Radiation condition)**

Let  $\omega$  be non-singular and  $\Pi_\pm$  the above projections. We say that  $u : \Omega \rightarrow \mathbb{C}$  with  $\|u\|_{sL} < \infty$  satisfies the radiation condition if

$$\Pi_-(u|_{W_r}) \rightarrow 0 \quad \text{and} \quad \Pi_+(u|_{W_{-r}}) \rightarrow 0 \quad \text{as } r \rightarrow +\infty$$

We identify a function on  $W_r$  with a function on  $W_0$  via a shift.

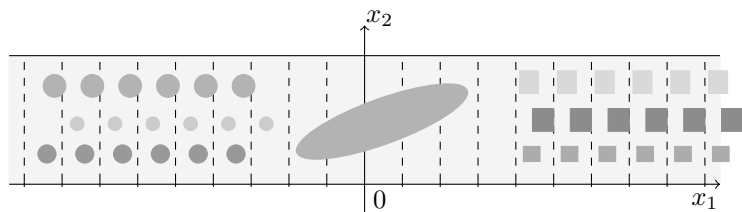
# Non-periodic media

## Theorem (Media that are periodic at infinity)

$\Omega = \mathbb{R} \times S$  be as above,  $a : \Omega \rightarrow \mathbb{R}^{d \times d}$  essentially bounded, symmetric and positive. Periodicity outside a compact set: For some  $R_0 > 0$

$$a(x + e_1) = a(x) \quad \text{for every } x \in \Omega \text{ with } |x_1| > R_0,$$

$\omega > 0$  non-singular frequency for the left and the right medium. If (i)–(iii) with  $f = 0$  possesses only the trivial solution, then there is a unique solution  $u$  for arbitrary  $f \in H^{-1}(\Omega)$  with compact support.



The theorem has the character of a Fredholm alternative

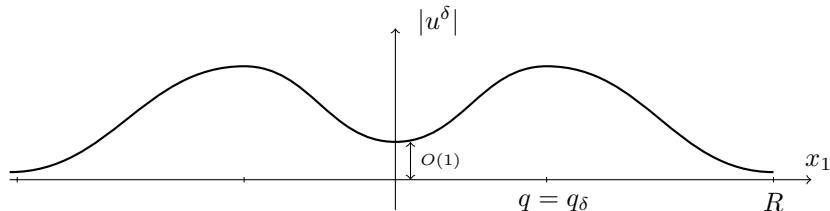
**Failure 1: Try Limiting absorption**

Solve on the infinite domain

$$-\nabla \cdot (a \nabla u^\delta) = (\omega^2 - i\delta)u^\delta + f$$

 $\delta > 0$  provides decay at infinity

Normalized in central region, norms may look like this:



$$\exists (q_\delta)_\delta : |u^\delta(q_\delta)| \rightarrow \infty$$

The limit  $\delta \rightarrow 0$  cannot be performed!



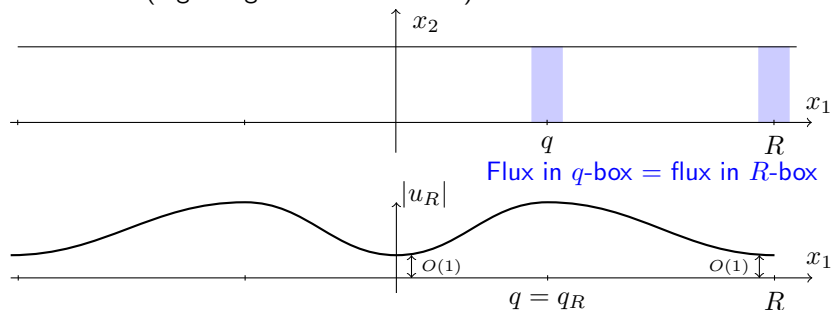
**Good idea:** Truncate to  $\Omega_R := \Omega \cap \{-R-1 < x_1 < R+1\}$  and seek

$$u = u_R \in V_R := \{u \in H^1(\Omega_R) \mid u|_{W_R} \in X_+, u|_{W_{-R-1}} \in X_-\}$$

which solves the Helmholtz equation between  $-R$  and  $R$ .

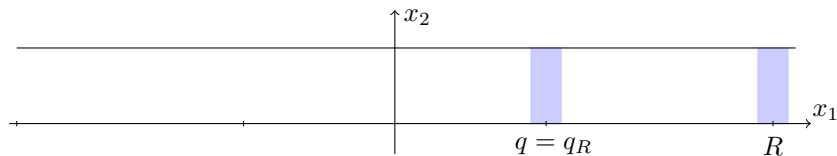
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**Failure 2** (regarding uniform estimates)



$u$  in  $q$  could be the sum of a large right-going and a large left-going wave!

## The successful proof



- In  $W_q$ :  $u = \Phi + \Psi + v$  with  $\Phi$  right-going,  $\Psi$  left-going,  $v$  small
  - Consider  $U := u - \Phi$ . Then:  $U$  is still right-going in  $W_R$ .  
And:  $U$  is approximately left-going in  $W_q$
  - Flux-equality for  $U$  implies that  $U$  is small in both boxes!
    - In  $W_R$ :  $u \approx \Phi$
    - In  $W_q$ :  $\Psi \approx 0$
- Hence:  $u$  is similar (in norms) in  $W_R$  and  $W_q$
- Contradiction argument, distinguishing cases regarding the positions  $q = q_R$  of “bad” boxes: Solution sequence remains bounded

We have shown: **existence result with energy methods**

- periodic media: existence and uniqueness
- piecewise periodic media: Fredholm alternative

Open questions

- limiting absorption principle
- characterization of non-singular frequencies
- beyond wave-guide geometries

# Thank you!